

Anhydrous Grotthuss mechanism for fast proton transport in a dense oxide-ion array of α -MoO₃

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Abstract

Developing high-power battery chemistry is an urgent task to buffer fluctuating renewable energies and achieve a sustainable and flexible power supply. Owing to the small size of proton and its ultrahigh mobility in water *via* the Grotthuss mechanism, aqueous proton batteries are an attractive candidate for high-power energy storage devices. Although Grotthuss proton transfer usually occurs in hydrogen-bonded networks of water molecules, in this work, we discover anhydrous Grotthuss-type proton transport in a dense oxide-ion array of solid α - MoO_3 even without structural water. The fast proton transfer and accumulation that occurs during (de)intercalation in α - MoO_3 is unveiled using both experiments and first-principles calculations. Coupled with a zinc anode and a superconcentrated dual-ion $\text{Zn}^{2+}/\text{H}^+$ electrolyte, the solid-state anhydrous Grotthuss proton transport mechanism realizes an aqueous MoO_3 -Zn battery with both high energy and power densities.

Main

The increasing concern given to the global environment and energy sustainability is driving the research and development of electrochemical energy storage devices that provide power supply with more resilience and flexibility. Currently, lithium-ion batteries dominate the power-source market for portable devices and electric vehicles due to their high energy density, high energy efficiency, and long lifetime.^{1,2} However, the global maldistribution of lithium resources has impeded their further widespread use.³ In particular, flammable organic electrolytes in lithium-ion batteries result in a low safety level and high fabrication/maintenance costs, both of which are unacceptable for grid-scale use.⁴ Therefore, aqueous rechargeable batteries that contain safe and less expensive aqueous electrolytes are an important future alternative for sustainable development.⁵

For the development of aqueous batteries with high energy density, exploiting Zn metal as a negative electrode is a straightforward approach because the Zn metal electrode possesses high theoretical capacities of 820 mAh/g and 5854 mAh/L. While reversible Zn metal plating/stripping is an important issue to be addressed for the development of aqueous Zn-ion batteries⁶, another large challenge also remains in a positive electrode. The intercalation of large hydrated Zn^{2+} generally causes damaging structural changes upon charge/discharge, leading to capacity degradation after cycling.⁷ Consequently, the concept of aqueous dual-ion batteries has recently been increasingly studied. For example, an aqueous $\text{Zn}^{2+}/\text{Li}^+$ dual-ion battery, which consists of a Zn metal plating/stripping negative electrode and Li^+ intercalation positive electrode (e.g., LiFePO_4) with an aqueous dual-ion $\text{Zn}^{2+}/\text{Li}^+$ electrolyte, was reported to provide an energy density of approximately 95 Wh/kg with 90% capacity retention after 80 cycles⁸.

In this work, we focus on protons as charge carriers in aqueous dual-ion batteries (**Figure 1**). Proton is the smallest and lightest cation; thus, it can be easily (de)intercalate in various structures at a fast rate.⁹ Moreover, barrierless H^+ hopping enables fast H^+ transport in an electrolyte owing to the Grotthuss

mechanism, where protons are transferred through the hydrogen bond network.¹⁰ In recent years, the Grotthuss topochemistry was extended to hydrate solid-state materials; for example, Prussian blue analogs exhibit fast H⁺ (de)intercalation with the assistance of structural water networks.¹¹ Similarly, the intercalated water layers in transition metal carbide nanosheets (MXenes) facilitate H⁺ storage and fast H⁺ transfer.^{12,13} Without structural/confined water, quinone-based organic compounds, which store H⁺ on carbonyl groups, exhibit both long cycle lives and large capacities.¹⁴

Among various transition metal (TM) oxides (TM= Mn, V, W, Ti, Mo)^{15–24} that deliver large capacities upon protonation, orthorhombic MoO₃ (α -MoO₃) possesses a unique bilayered structure (**Figure 1**) that accommodates various cations, such as Li⁺^{25,26}, Na⁺²⁷, Ca²⁺²⁸ and Mg²⁺²⁹ in organic electrolytes. In aqueous electrolytes, the intercalation of bare Zn²⁺^{30,31} and Al³⁺³² has also been studied. However, the H⁺ intercalation behavior in α -MoO₃ with various aqueous electrolytes remains controversial: some reports claim a bare H⁺ intercalation mechanism^{33–37} that is consistent with H_xMoO₃ bronze obtained through the spillover method³⁸, while others report water and H⁺ cointercalation^{39,40}. Meanwhile, α -MoO₃ suffers severe dissolution in aqueous electrolytes; therefore, the capacity decays gradually upon cycling, making it difficult to study the detailed H⁺ intercalation mechanism. Although many attempts have been made to solve this issue, for example, the use of highly concentrated electrolytes^{40,41}, gel-type or quasi-solid-state electrolytes with polymer additives³¹, and electrode surface coating with polymers or ceramics^{30,42}, the detailed H⁺ intercalation mechanism in α -MoO₃ has yet to be fully understood. Herein, we provide the full comprehension of H⁺ intercalation in α -MoO₃ as a cathode material for aqueous Zn²⁺/H⁺ batteries. Fast H⁺ transfer in α -MoO₃ through a solid-state anhydrous Grotthuss mechanism realizes aqueous batteries with both high power and high energy densities.

Electrochemical properties of MoO₃ in aqueous Zn²⁺/H⁺ electrolytes

α -MoO₃ was synthesized by a previously reported hydrothermal method,⁴⁰ and the synchrotron X-ray diffraction pattern confirmed the successful synthesis of a pure α -MoO₃ phase (**Figure S1**). To study the H⁺ intercalation mechanism in α -MoO₃, ZnCl₂ was selected as an electrolyte salt. In addition to its compatibility with a Zn anode, the high solubility of ZnCl₂ enabled the formation of a superconcentrated liquid structure with limited amount of free water molecules, while its Brønsted acidity generated a low pH environment with a high H⁺ concentration. Therefore, the electrochemical properties of α -MoO₃ were evaluated using three aqueous electrolytes: conventional Zn²⁺ electrolyte (3 mol kg⁻¹ ZnCl₂/H₂O), superconcentrated Zn²⁺ electrolyte (32 mol kg⁻¹ ZnCl₂/H₂O), and superconcentrated dual-ion (Zn²⁺/H⁺) electrolytes (32 mol kg⁻¹ ZnCl₂ + 1 mol kg⁻¹ P₂O₅/H₂O). Note that P₂O₅ in a dual-ion electrolyte

generates H^+ through hydrolysis ($\text{P}_2\text{O}_5 + 3\text{H}_2\text{O} \rightarrow 2\text{H}_3\text{PO}_4$). The Raman spectra for the superconcentrated electrolytes indicate that most water molecules are coordinated to Zn^{2+} (**Figure S2**).

Before testing the electrochemical properties of $\alpha\text{-MoO}_3$, we evaluated the negative electrode, namely, Zn stripping and plating on a Ti current collector using the three electrolytes (**Figures S3 and 2a**).⁶ The average coulombic efficiency in the aqueous $\text{Zn}^{2+}/\text{H}^+$ dual-ion electrolyte is 99.0% over 200 cycles (**Figure S3**), largely outperforming the superconcentrated Zn^{2+} electrolyte (95.0% over 200 cycles) and conventional Zn^{2+} electrolyte (82.2% over 100 cycles). The improved zinc reversibility after the addition of P_2O_5 may result from the formation of a $\text{Zn}_3(\text{PO}_4)_2$ -based solid electrolyte interphase (SEI) layer.⁴³ This reversible Zn stripping and plating was used as the counter electrode in this work.

Figure 2b shows the cyclic voltammetry (CV) curves measured for $\alpha\text{-MoO}_3$ using the three electrolytes at a scan rate of 0.5 mV s^{-1} . To suppress the hydrogen evolution and chlorine evolution reactions on the $\alpha\text{-MoO}_3$ electrode, the cutoff voltages were set at 0.45 and 1.3 V for the aqueous $\text{Zn}^{2+}/\text{H}^+$ dual-ion electrolyte (**Figure 2a**) and 0.25 and 1.1 V for the Zn^{2+} aqueous electrolytes. In the first CV cycle (**Figure 2b inset**), the $\alpha\text{-MoO}_3$ electrode shows identical asymmetric-shaped CV curves for all three electrolytes, in which there are four cathodic and two anodic current flows. However, in the subsequent cycles, the CV curves become symmetric, showing two pairs of redox peaks. Note that the redox potentials using the aqueous $\text{Zn}^{2+}/\text{H}^+$ dual-ion electrolyte are centered at approximately 0.6 and 0.95 V vs. Zn/Zn^{2+} , which shift by +0.2 V from those using Zn^{2+} aqueous electrolytes presumably owing to the potential shift of the Zn/Zn^{2+} counter electrode and/or the change in the activity of H^+ . The shapes of the CV curves for all three electrolytes resemble those reported for H^+ (de)intercalation in $\alpha\text{-MoO}_3$ using a 9.5 M H_3PO_4 aqueous electrolyte⁴⁰, suggesting dominant H^+ (de)intercalation even when using the Zn^{2+} aqueous electrolytes. Indeed, the X-ray fluorescence (XRF) elemental analysis of the electrodes after a cathodic scan shows no evident increase in the peak intensity of Zn compared to that of the pristine electrode (**Figure. S4**). Considering that the aqua Zn^{2+} complex is a Brønsted acid to generate H^+ , it is most likely that H^+ (de)intercalation occurs even when using the aqueous Zn^{2+} electrolytes. Importantly, while 3 and 32 mol kg^{-1} ZnCl_2 aqueous electrolytes exhibit steep redox peak degradation upon repeated CV cycling, the 32 mol kg^{-1} $\text{ZnCl}_2 + 1 \text{ mol kg}^{-1}$ P_2O_5 aqueous dual-ion electrolyte exhibits stable CV curves (**Figure 2b**). The improved cycle stability should be ascribed to the suppression of $\alpha\text{-MoO}_3$ dissolution and the formation of effective SEI.⁴³ Indeed, the X-ray photoelectron spectroscopy (XPS) analysis of the Zn metal anode after cycling in the 3 m ZnCl_2 electrolyte evidences the Mo deposition from the Mo ions dissolved in the electrolyte (**Figure S6**).

Figure 2c shows the charge/discharge curves of the $\alpha\text{-MoO}_3$ electrode with galvanostatic charging followed by 3 h of potentiostatic charging in the aqueous $\text{Zn}^{2+}/\text{H}^+$ dual-ion electrolyte. The $\alpha\text{-MoO}_3$

electrode delivers a large capacity of 465 mAh g⁻¹ at a rate of 0.5 A g⁻¹ during the first discharge, corresponding to 2.5 H⁺ intercalation per formula unit of MoO₃ with an average voltage of approximately 0.9 V. Note that 'discharge' and 'charge' of the α-MoO₃ electrode are defined as H⁺ intercalation (cathodic process) and deintercalation (anodic process), respectively. Although the galvanostatic charge at 0.5 A g⁻¹ can extract only 1.5 H⁺, the remaining 1.0 H⁺ can be extracted when applying a constant voltage of 1.3 V for 3 h (**Figures S7** and **S8**). The diffusion coefficient determined by the potentiostatic intermittent titration technique (PITT) shows a significant 4.9-fold deceleration in H⁺ diffusion during the deprotonation from H_{1.1}MoO₃ to MoO₃, confirming the trapped nature of ~1.0 H⁺ in MoO₃ (**Figure S9**).

Under galvanostatic charging (without a potentiostatic step), the α-MoO₃ electrode in an aqueous Zn²⁺/H⁺ dual-ion electrolyte retains 98% of its initial capacity after 1000 cycles at a rate of 2 A g⁻¹ (**Figure 3a**). Furthermore, 62% of the specific capacity at 1 A g⁻¹ is available at the fast discharge rate of 16 A g⁻¹ (**Figure 3b** and **Figure S10**). These performance results indicate the stability of the MoO₃ framework against (de)protonation as well as the fast proton diffusion therein H_xMoO₃ (1.0 ≤ x ≤ 2.5). In contrast, the α-MoO₃ electrodes in the aqueous Zn²⁺ electrolytes have capacity retentions of only 24.5% and 63.8%, respectively (**Figures S11** and **S12**). Moreover, both the capacity and cycling stability of the α-MoO₃ electrode in the aqueous Zn²⁺/H⁺ dual-ion electrolyte outperform those reported previously for α-MoO₃ electrodes using aqueous electrolytes, including quasi-solid-state Zn²⁺ electrolytes^{30,31} and concentrated acid electrolytes^{33,40}.

MoO₃ host-lattice response to proton intercalation

To clarify the structural evolution of the α-MoO₃ electrode in the aqueous Zn²⁺/H⁺ dual-ion electrolyte, *in situ* X-ray diffraction (XRD) was performed during the 1st cycle (**Figure 4a**). The interlayer distance of α-MoO₃ remains nearly constant (approximately 7.0 Å) during the entire protonation process, while it increases from 7.0 to 7.5 Å and then decreases to 7.0 Å during deprotonation. This asymmetric lattice response is consistent with the asymmetric CV and charge/discharge curves, which also highly resemble those reported in a 6 M H₂SO₄ electrolyte³³ and in a 4.4 M H₂SO₄ electrolyte³⁵. However, despite the asymmetric structural evolution, the α-MoO₃ structure recovers to the pristine state after a constant voltage is applied. The structure of the fully protonated phase was clarified using *ex situ* synchrotron XRD and Rietveld refinement (**Figure 4b** and **4c**). Although it is difficult to determine proton positions using X-rays, the MoO₃ framework only exhibits a slight monoclinic distortion after protonation, which is consistent with a previous report on H_{1.68}MoO₃⁴⁴ and the *in situ* XRD results.

Importantly, no water cointercalation occurs in the protonated structure. When a small amount of water (one water or hydronium intercalation in 16 formula units of MoO₃) is intercalated, the density functional theory (DFT) calculation predicts that the interlayer distance expands by 13.8% (H₂O intercalation) and

17.2% (H_3O^+ intercalation) (**Figure S13**), which are considerably larger than those observed experimentally. Additionally, the interlayer distance after water cointercalation in a 1 M H_2SO_4 electrolyte has been reported to show an expansion of 11% upon protonation³⁹, in contrast to the negligible change observed in the protonation process in our experiment. Therefore, bare H^+ (de)intercalation occurs in the $\alpha\text{-MoO}_3$ electrode with the aqueous $\text{Zn}^{2+}/\text{H}^+$ electrolyte.

Solid-state anhydrous Grotthuss mechanism

The above experimental results indicate asymmetric bare- H^+ (de)intercalation in the $\alpha\text{-MoO}_3$ electrode. To clarify the origin of this asymmetry, we conducted DFT calculations on the H^+ dynamics in $\alpha\text{-MoO}_3$. The stable sites for H^+ were determined by the structural optimization of the protonated phases. The three most stable H^+ absorption sites are labeled as site A ($\text{O2-H}\cdots\text{O2}$), site B ($\text{O1-H}\cdots\text{O1/O2}$) and site C ($\text{O2-H}\cdots\text{O1}$), where O1 and O2 are the terminating oxygen coordinated to one Mo and the bridging oxygen coordinated to two Mo, respectively (**Figure S14**). The remaining edge-sharing oxygen (O3) coordinated to three Mo atoms is not favorable for H^+ adsorption because all O $2p$ orbitals ($2p_x$, $2p_y$, and $2p_z$) participate in the Mo $4d\text{-O } 2p$ bonds. Site A is located within the MoO_3 layer, while sites B and C are located in the MoO_3 interlayer space.

The most favorable site for H^+ changes as a function of x in H_xMoO_3 . At low H^+ concentrations during protonation ($0 < x \leq 0.5$ in H_xMoO_3), intralayer site A is the most favorable for H^+ adsorption. However, after 0.5 H^+ intercalation ($0.5 < x \leq 2.5$ in H_xMoO_3), interlayer sites B and then C become more favorable than site A, leading to the final formula of $\text{H}_{2.5}\text{MoO}_3$ (site A (0.5H^+) \rightarrow site B (1.5H^+) \rightarrow site C (0.5H^+)), **Figure 5a** and **Figure S15**). Further protonation of site B (H_3MoO_3) is not thermodynamically favored, as it causes structural decomposition.

On the other hand, at the early stage of deprotonation ($2.5 > x \geq 2.25$ in H_xMoO_3), the DFT calculations suggest that 0.25 H^+ is extracted from interlayer site C. Then, for $2.25 > x \geq 1.5$ in H_xMoO_3 , 0.5 H^+ at site A is predominantly deintercalated in addition to the simultaneous partial deintercalation of H^+ at site B. After completing H^+ deintercalation from site A, the H^+ at site B is deintercalated ($1.5 > x \geq 0$ in H_xMoO_3). In parallel, the remaining 0.25 H^+ at site C is deintercalated at the end of deprotonation. Notably, the deprotonation order (site C (0.25H^+) \rightarrow site A (0.5H^+) \rightarrow site B (0.75H^+) \rightarrow site B + site C (1.0H^+)) (**Figure 5b**) does not follow the protonation order (site A (0.5H^+) \rightarrow site B (1.5H^+) \rightarrow site C (0.5H^+)), which explains the asymmetric charge/discharge profile. Experimentally, complete deprotonation is possible only when a constant voltage is applied (**Figure 2c**). The sluggish deprotonation at the end of charging may arise from the slow kinetics of the remaining H^+ at sites B and C. The calculated voltage profile agrees well with the experimental results, confirming the validity of the asymmetric (de)protonation processes (**Figure S16**).

To unveil the origin of the facile (de)protonation in the $1.0 < x \leq 2.5$ range of $H_x\text{MoO}_3$ and the sluggish deprotonation in the $1.0 \geq x \geq 0$ range of $H_x\text{MoO}_3$, climbing image nudged elastic band (CI-NEB) calculations were conducted, and the energy profile of H^+ diffusion in MoO_3 was visualized (**Figure 5c**). At the beginning of protonation, H^+ rotates and hops consecutively between site A with a low energy barrier of 0.13 eV. The dense, zigzagged O2 (bridging oxygen) array provides a 1D channel for fast proton transport within the MoO_3 intralayer space, where the short O2-O2 distance (2.63 Å) facilitates proton hopping. Upon the further protonation of $H_{0.5}\text{MoO}_3$, the interlayer diffusion channel containing site B also has a low energy barrier of 0.26 eV with an O1-O1/O2 distance of 2.71 Å (**Figure S17**), confirming the high mobility of H^+ at site B.

As shown in **Figure 5d**, the deprotonation from site C of $H_{2.5}\text{MoO}_3$ also involves a small energy barrier of 0.29 eV through the interlayer diffusion channel containing sites B and C (O-O distance of 2.86 Å). The H^+ transfer in the dense oxide-ion array is described as the solid-state anhydrous Grotthuss mechanism, whose kinetics highly rely on the distance between two adjacent lattice oxide ions. **Figure S18** summarizes the statistics of O-O distances in all possible deprotonation channels of the charged structures. As expected, during deprotonation from $H_{2.5}\text{MoO}_3$ to HMoO_3 , short O-O distances of < 2.9 Å exist to enable solid-state anhydrous Grotthuss H^+ transfer. However, after extracting 0.25 H^+ from site C in $H_{2.5}\text{MoO}_3$, the O-O distances in the diffusion channel consisting of sites B and C increase drastically (> 3.1 Å); meanwhile, the energy barrier therein for H^+ diffusion becomes remarkably high (0.98 eV, **Figure S19**). Therefore, further deprotonation from site C is unfavorable, instead, protons in site A are deintercalated, resulting in asymmetric (de)protonation (**Figure 5b**). Indeed, after deprotonation to HMoO_3 , all the long-range 1D channels are disrupted so that fast Grotthuss H^+ transfer is no longer applicable (**Figure S20**). The remaining H^+ can be removed only under long relaxation times, such as potentiostatic charging and PITT (**Figure S21** and **S22**). However, except for the trapped H^+ at site C, $H_x\text{MoO}_3$ ($1.0 \leq x \leq 2.5$) exhibits fast H^+ transport through the diffusion channels upon charge/discharge, providing a remarkably high capacity and high-rate capability, as demonstrated in **Figure 3**.

To summarize, coupled with a zinc metal anode and an aqueous dual-ion $\text{Zn}^{2+}/\text{H}^+$ (32 m ZnCl_2 + 1 m P_2O_5) electrolyte, the MoO_3 -Zn battery delivers a large energy density of 413 Wh kg^{-1} upon discharge at a power density of 0.90 kW kg^{-1} as well as a peak power density of 10.52 kW kg^{-1} at an energy density of 217 Wh kg^{-1} per weight of MoO_3 ; these results are more than double that of a similar MoO_3 -Zn battery with a ZnCl_2 -based electrolyte³⁹ (energy density of 198 Wh kg^{-1} at a power density of 0.28 kW kg^{-1} and power density of 6.7 kW kg^{-1} at an energy density of 104.5 Wh kg^{-1}). Moreover, with the aid of the solid-state anhydrous Grotthuss mechanism, this prototype cell successfully outperforms most aqueous zinc-ion batteries and proton batteries (**Figure S23**). Contrary to conventional intercalation chemistry, which requires a porous host that accommodates ion diffusion and storage, the solid-state anhydrous Grotthuss mechanism demonstrated in this work enables fast H^+ transfer and accumulated H^+ storage in dense oxide-ion arrays. Therefore, further exploration for other host materials capable of H^+

intercalation based on the solid-state anhydrous Grotthuss mechanism will be an important challenge for not only fabrications of high-power aqueous H⁺ batteries but also other solid-state ionics applications using protons.

Methods

α -MoO₃ was synthesized *via* a simple hydrothermal synthesis approach⁴⁰, where 1 g of (NH₄)₆Mo₇O₂₄·4H₂O (Wako) was dissolved in 25 mL of water, followed by the addition of 10 mL of HNO₃ solution (3 M). After stirring for 10 min, the transparent colorless solution was poured into a 50 mL TeflonTM-lined Parr autoclave and heated at 180 °C for 12 h. After filtration, the white-colored powder was washed with water and ethanol and then dried in an oven at 80 °C in air overnight.

The electrolyte was prepared by weighing ZnCl₂ and P₂O₅ in an argon-filled glovebox to prevent the absorption of water. After weighing, ultrapure water (Wako) was mixed with salts to form a solution in open air.

The working electrode for electrochemical characterization was formulated by slurry casting 70 wt% MoO₃ active material, 20 wt% Ketjen-black (ECP-600JD, Lion Corp.) and 10 wt% polyvinylidene difluoride (PVDF) binder on Ti foil or carbon paper (for *in situ* X-ray diffraction (XRD)) using prompt amounts of N-methyl-2-pyrrolidone (NMP) (Kanto, 99%) as the solvent. The loading level of the electrodes was controlled to ~1 mg cm⁻².

The electrochemical performance was tested in PTFE three-electrode beaker cells with Ag/AgCl (in a saturated KCl aqueous solution) as the reference electrode and activated carbon as the counter electrode. In the two-electrode cell tests, zinc foil was attached to a titanium wire, working as both the reference and counter electrode. The zinc stripping and plating tests were conducted in a 2032-type coin cell where a titanium foil and a zinc foil were separated with a glass fiber separator (Fisher) presoaked with an electrolyte. The *in situ* XRD studies were conducted in a custom cell with a Kapton[®] membrane window on the cathode side.

The electrochemical performance was studied *via* cyclic voltammetry (CV) and galvanostatic charge and discharge (GCD), the galvanostatic intermittent titration technique (GITT) and the potentiostatic intermittent titration technique (PITT) using a VMP3 potentiostat (BioLogic) at room temperature.

Powder XRD and *in situ* XRD studies were conducted with a Bruker-AXS D8 ADVANCE diffractometer using a Co sealed tube (operating at 35 kV, 40 mA) and a linear position-sensitive detector (LYNX-EYE). VESTA⁴⁵ software was used to illustrate the crystal structure. For crystal structure refinement, we used a powder diffractometer at a synchrotron radiation beamline 5S2 of Aichi Synchrotron Radiation Center,

Japan. The wavelength used was calibrated by refining a powder diffraction pattern of the ceria power, and the value was 0.700072(9) Å.

Elemental analysis of the MoO₃ electrodes by energy dispersive X-ray fluorescence (XRF) spectroscopy was performed with a JSX-3400RII instrument (JEOL).

The surface chemistry of the Zn electrodes was analyzed *via* X-ray photoelectron spectroscopy (XPS, PHI5000 VersaProbe II, ULVAC-PHI) with a monochromatic Al K_X-ray source. The Zn electrode samples were rinsed with pure water and acetone after being extracted from cycled cells.

The coordination states of the water molecules were investigated by Raman spectroscopy (NRS-5100, JASCO). A 532 nm excitation laser was used. The electrolyte solution was sealed in a quartz cell, and the laser was directed through the quartz crystal window.

Spin-polarized DFT calculations were carried out using the Vienna ab initio simulation package (VASP).^{46,47} The projector-augmented wave pseudopotential and a plane-wave basis set with an energy cutoff of 520 eV were used.⁴⁸ The Perdew-Burke-Ernzerhof functional^{49,50} with the Grimme scheme D2⁵¹ described the exchange-correlation energy, and the Hubbard U correction was applied to the d electrons of the Mo atoms ($U_{\text{eff}} = 3.5$ eV).⁵² A k-point separation of approximately 0.04 Å⁻¹ was used in all calculations.

Based on a 2 × 1 × 2 supercell (Mo₁₆O₄₈), we calculated a series of H-intercalated structures of H_xMoO₃ (0 ≤ x ≤ 3.0). Each proton was placed between two nearby O atoms where one covalent bond (~1 Å) and one hydrogen bond (1.6~2.3 Å) were formed. All possible configurations were calculated by optimizing both the lattice and ion positions, and then we constructed a convex hull to identify the stable phases in the discharge and charge processes. The formation energies (at 0 K) were calculated as:

$$\Delta E_f(\text{H}_x\text{MoO}_3) = E(\text{H}_x\text{MoO}_3) - \frac{(3-x)}{3}E(\text{MoO}_3) - \frac{x}{3}E(\text{H}_3\text{MoO}_3) \quad (1)$$

where $E(\text{H}_x\text{MoO}_3)$, $E(\text{MoO}_3)$ and $E(\text{H}_3\text{MoO}_3)$ are the total energies of H_xMoO₃, pristine MoO₃ and H₃MoO₃ (maximally H-intercalated configuration), respectively. The determined structures are shown in Supplementary **Figure S19**. The voltage profile was then evaluated, with the average voltage of the reaction between two adjacent stable phases calculated as:

$$U(x_1, x_2) = - [E(\text{H}_{x_1}\text{MoO}_3) - E(\text{H}_{x_2}\text{MoO}_3) - (x_1 - x_2)\mu_{\text{H}}] / (x_1 - x_2)F \quad (2)$$

where F is the Faraday constant, μ_{H} is the chemical potential of hydrogen gas at atmospheric pressure, and T = 298.15 K.

$$\mu_{\text{H}} = E(\text{H}_2) + 7/2 k_{\text{B}}T - T\Delta S \quad (3)$$

In the above equation, $E(\text{H}_2)$ is the calculated total energy of a hydrogen molecule, k_{B} is the Boltzmann constant and ΔS is obtained from the JANAF thermochemistry tables.⁵³

We used the climbing image nudged elastic band (CI-NEB) method to calculate proton diffusion pathways in a $2\sqrt{2} \times 1 \times 2\sqrt{2}$ supercell.⁵⁴ For the stable structures in the discharge process, an extra proton was added to migrate between two neighboring adsorption sites. In the charge process, a proton was removed from the stable structures to construct a vacancy site for the migration of other protons.

Declarations

Author Contributions

M.O. and A.Y. conceived and directed the project. Z.M. conducted material synthesis and electrochemical measurements. Z.M. and X.M.S. conducted the theoretical calculations. S.N. conducted the structural analysis. Z.M. and S.K. designed the experiments. All authors wrote the manuscript.

Conflicts of Interest

There are no conflicts of interest.

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Figures

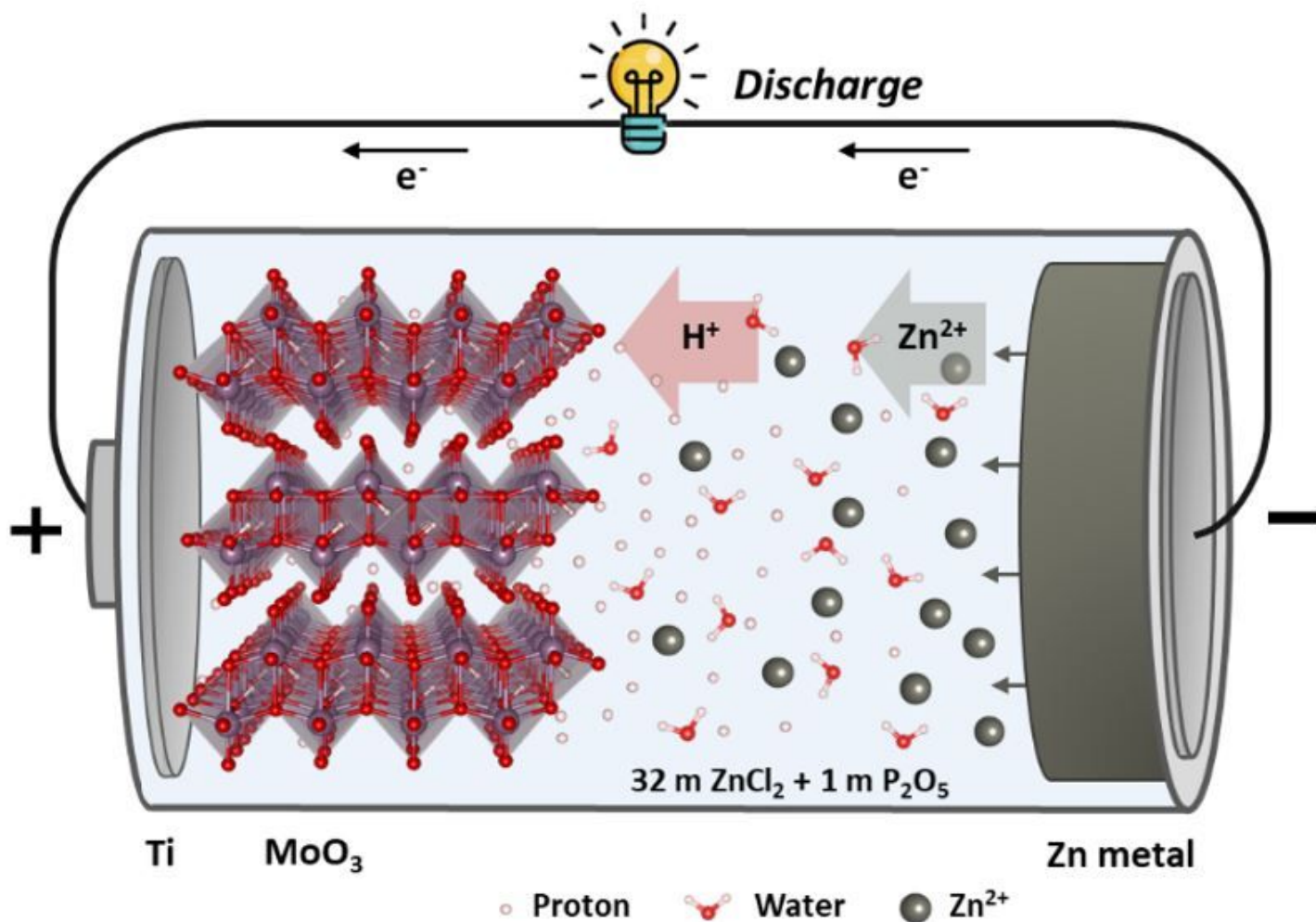


Figure 1

Schematic illustration of the aqueous dual-ion Zn²⁺/H⁺ battery prototype. The battery consists of a H⁺ intercalation positive electrode α -MoO₃ and a Zn metal plating/stripping negative electrode with an aqueous dual-ion Zn²⁺/H⁺ electrolyte.

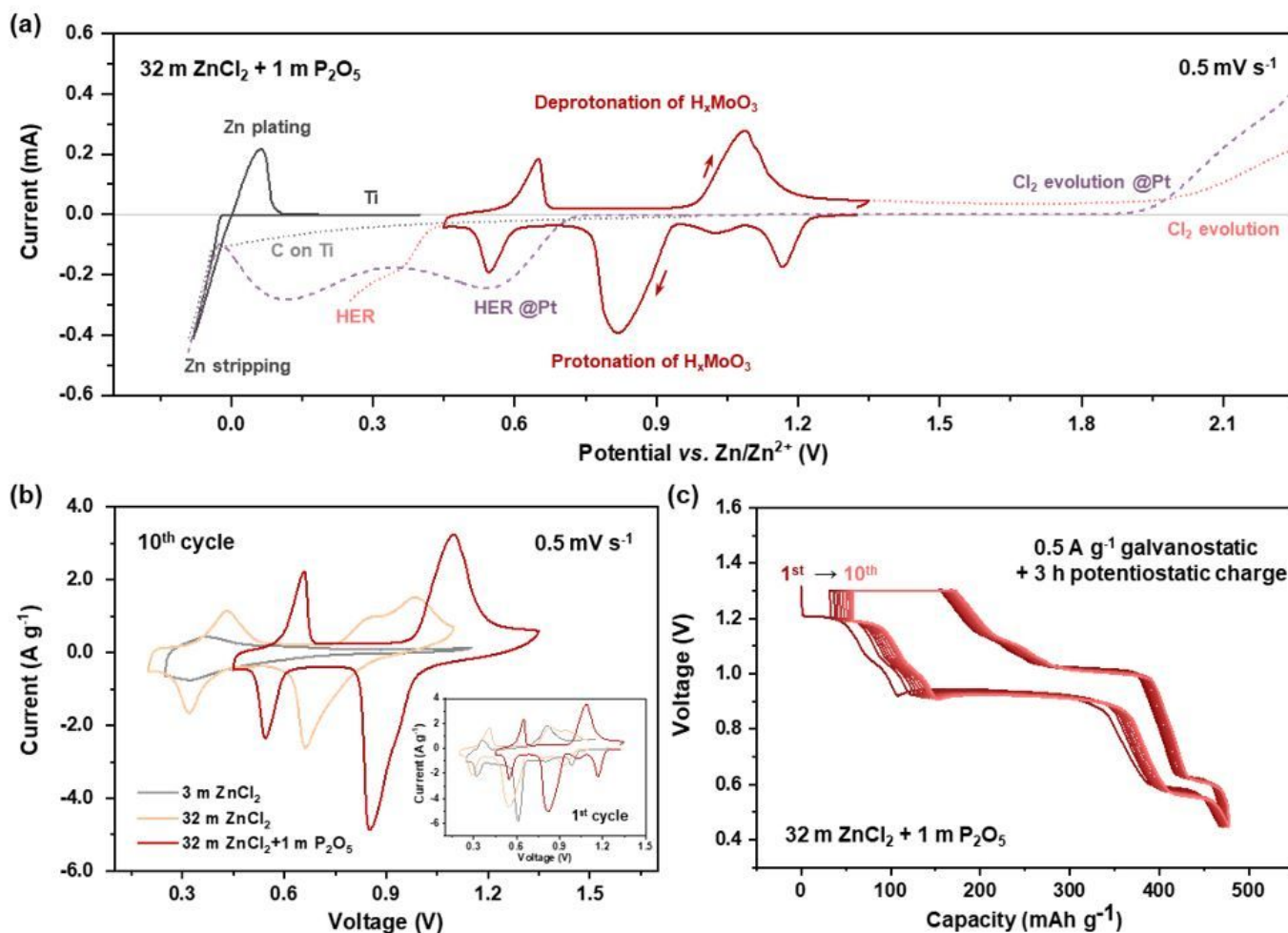


Figure 2

Electrochemical properties of α - MoO_3 in aqueous dual-ion $\text{Zn}^{2+}/\text{H}^+$ electrolytes. (a) CV curves of α - MoO_3 in a $32 \text{ m ZnCl}_2 + 1 \text{ m P}_2\text{O}_5$ aqueous electrolyte (red solid line). The CV curve of Zn plating/stripping on a Ti current collector (gray solid line) and the LSV curves of hydrogen-evolution and chlorine evolution reactions on a Pt current collector are also shown for comparison (purple dash lines). The red dotted lines are the LSV curves using the α - MoO_3 electrode. The gray dotted line is the LSV curve using Ti foil coated with carbon. (b) The CV curves of α - MoO_3 during the 10th cycle using 3 m ZnCl_2 (gray line), 32 m ZnCl_2 (orange line), and $32 \text{ m ZnCl}_2 + 1 \text{ m P}_2\text{O}_5$ (red line) aqueous electrolytes. The inset shows the CV curves at the first cycle. (c) The charge/discharge curves of α - MoO_3 in a $32 \text{ m ZnCl}_2 + 1 \text{ m P}_2\text{O}_5$ aqueous electrolyte under a galvanostatic + potentiostatic charging mode (A 3 h constant voltage of 1.3 V vs. Zn/Zn^{2+} was applied after each constant current charging at 0.5 A/g), corresponding to H_xMoO_3 ($0 \leq x \leq 2.5$).

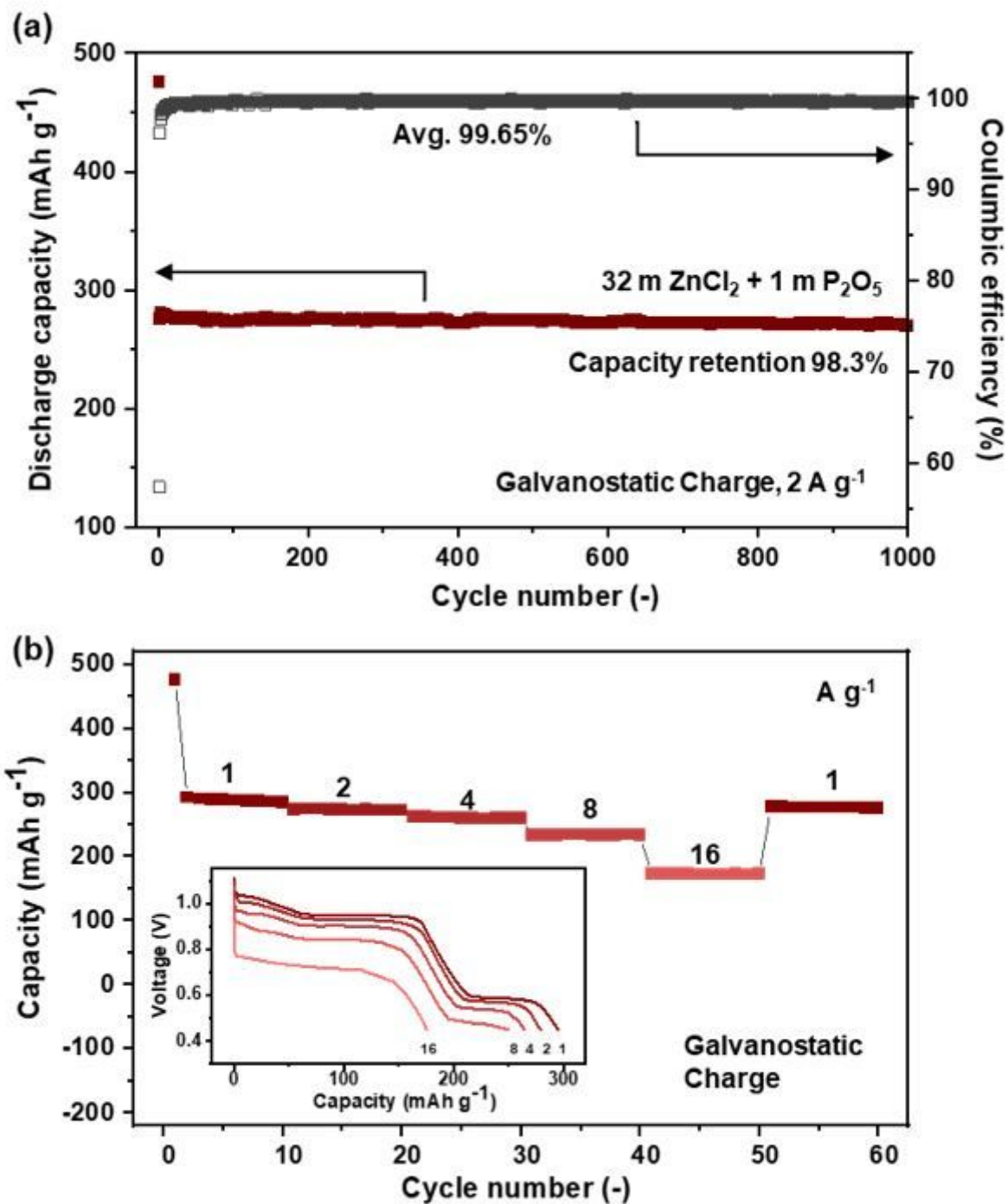


Figure 3

Cycle stability and rate capability of H_xMoO_3 ($1.0 \leq x \leq 2.5$) in Galvanostatic charge mode. (a)

Discharge capacity and coulombic efficiency at 2 A g⁻¹, (b) rate performance at discharge rates from 1 A g⁻¹ to 16 A g⁻¹. Inset: the corresponding discharge curves.

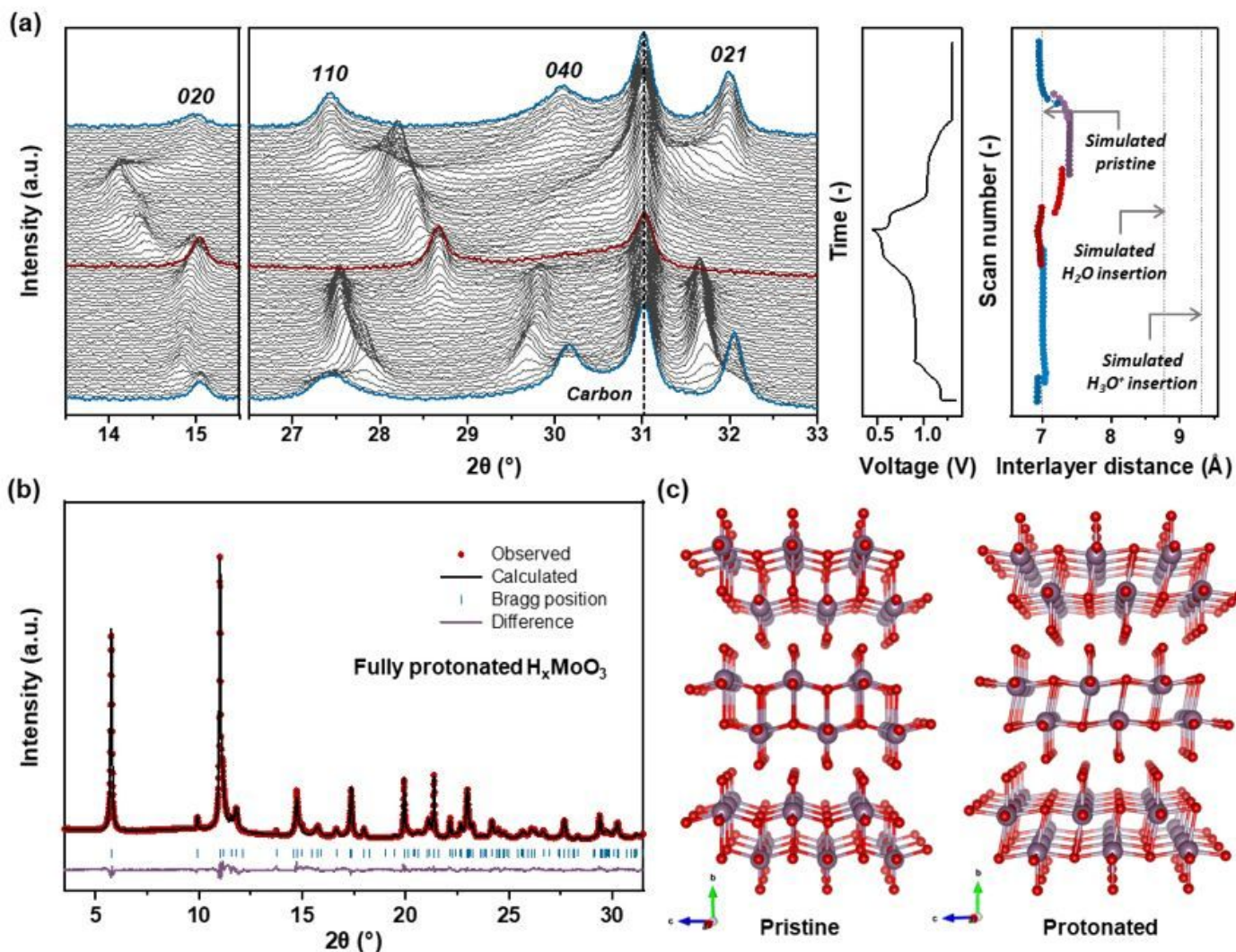


Figure 4

Structural evolution of MoO_3 upon (de)protonation. (a) *In situ* XRD patterns upon protonation and deprotonation of MoO_3 , and deduced interlayer-distance change (colored dotted lines). The simulated interlayer distances of H_2O (1 H_2O in 16 MoO_3) or H_3O^+ (1 H_3O^+ in 16 MoO_3) intercalated MoO_3 are also indicated for comparison (gray dotted lines). (b) Synchrotron XRD pattern and the Rietveld refinement result of protonated $\text{H}_{2.5}\text{MoO}_3$. (c) Refined structures of pristine (left) and protonated (right) phases. Note that the proton positions cannot be determined by the synchrotron XRD and therefore are not shown.

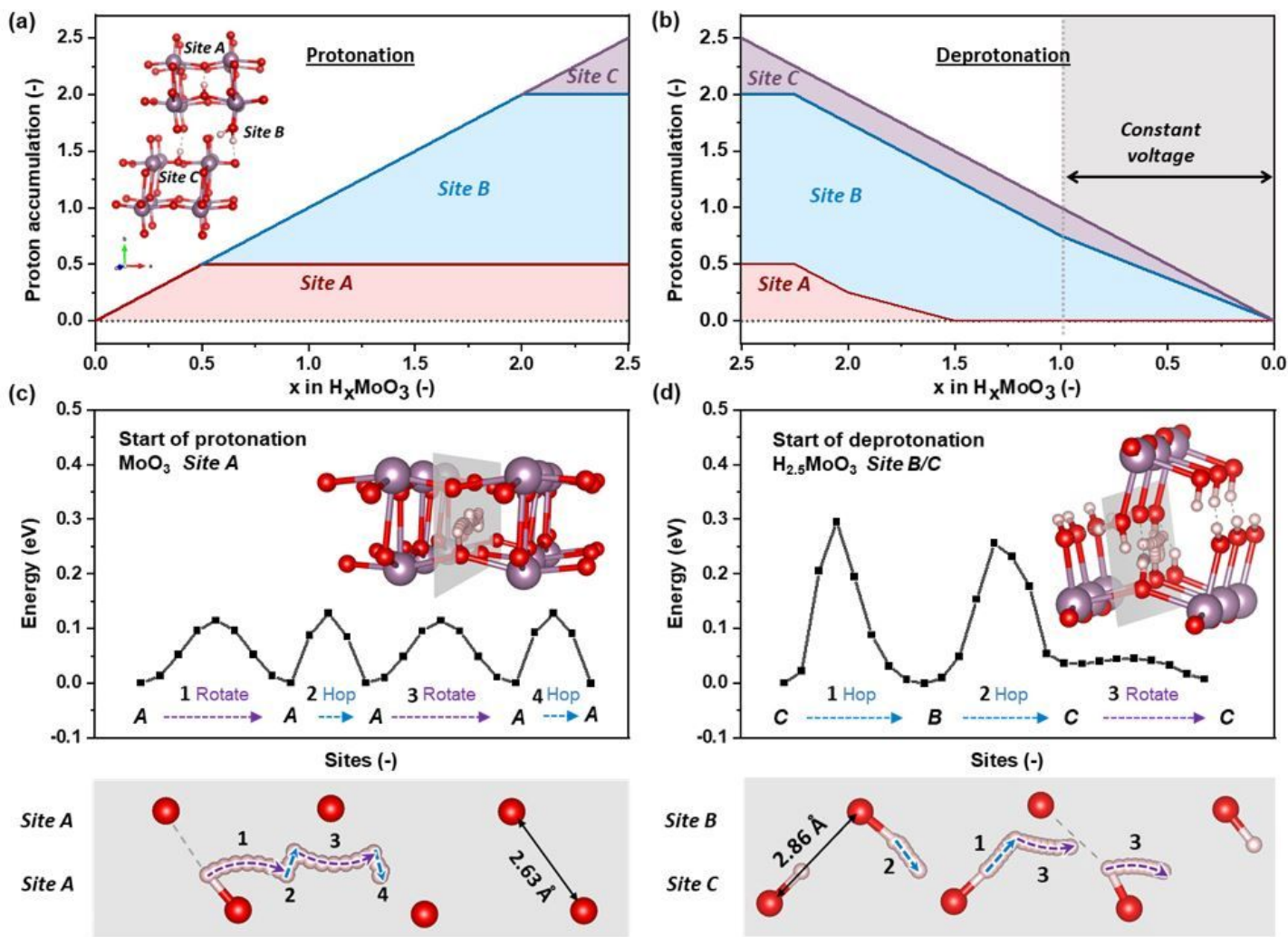


Figure 5

Solid-state anhydrous Grotthuss mechanism for fast proton transport in dense oxide-ion array. The occupancy of each site in (a) protonation with an order of site A ($0.5 H^+$) \rightarrow site B ($1.5 H^+$) \rightarrow site C ($0.5 H^+$), and (b) deprotonation with an order of site C ($0.25 H^+$) \rightarrow site A ($0.5 H^+$) \rightarrow site B ($0.75 H^+$) \rightarrow site B + site C ($1.0 H^+$). (c-d) Anhydrous Grotthuss mechanism for proton transport and corresponding energy barriers at the early stages of protonation (c) and deprotonation (d).

Supplementary Files

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